

Quantum Theory of Superresolution for Incoherent Optical Imaging *

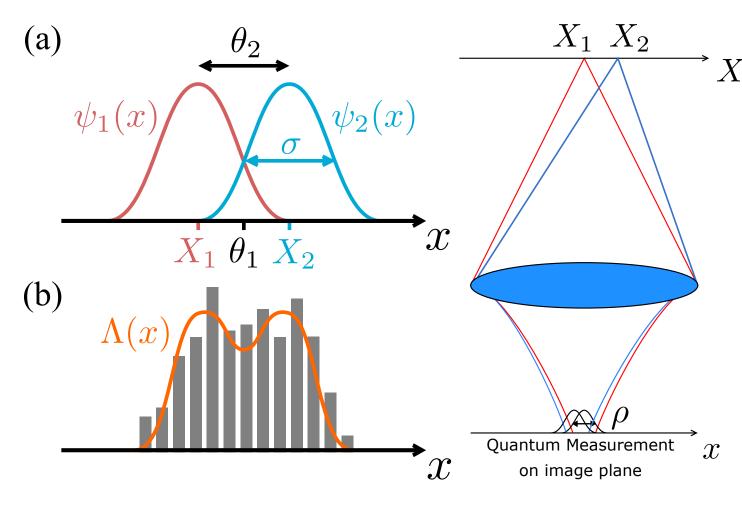
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March 2017

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Superresolution Incoherent Imaging



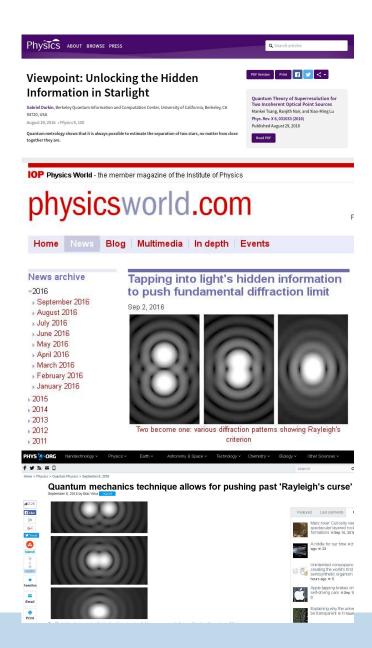
- 1. Tsang, Nair, and Lu, PRX **6**, 031033 (2016).
- Nair and Tsang, Opt. Express 24, 3684 (2016).
- Tsang, Nair, and Lu, Proc. SPIE 10029, 1002903 (2016).
- 4. Nair and Tsang, PRL (Editors' Suggestion) 117, 190801 (2016).
- 5. Tsang, arXiv:1605.03799 (2016).
- 6. Ang, Nair, and Tsang, arXiv:1606.00603 (2016).
- 7. Tsang, arXiv:1608.03211 (2016).
- 8. Lu, Nair, Tsang, arXiv:1609.03025 (2016).

Popular Coverage

- Viewpoint in APS Physics
- IoP Physics World and nanotechweb.org:

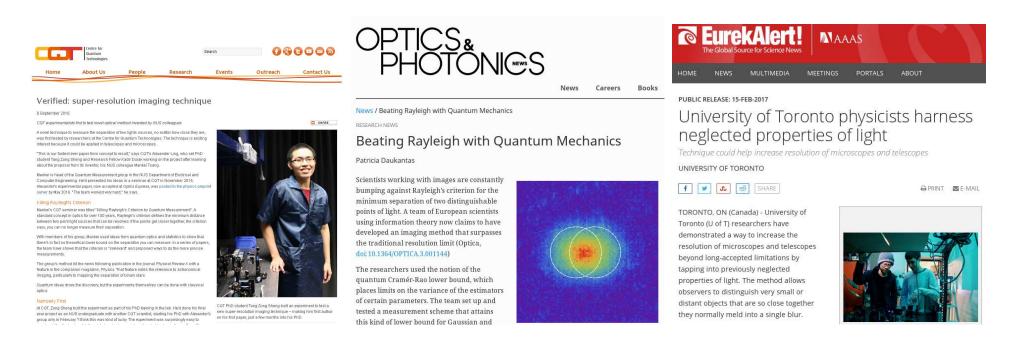
Seth Lloyd of the Massachusetts Institute of Technology in the US is impressed. 'This is awesome work and I am amazed that it hasn't been done before,' he says. 'Perhaps everyone thought it was too good to be true.'

- APS Physics Central
- Phys.org

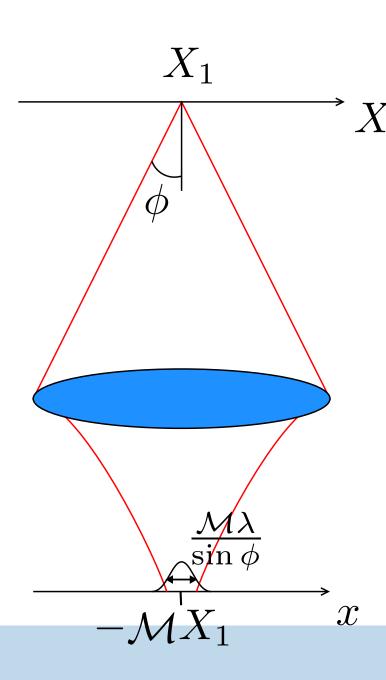


Experiments

- Tang, Durak, and Ling, "Fault-tolerant and finite-error localization for point emitters within the diffraction limit," Optics Express **24**, 22004 (2016).
- Yang, Taschilina, Moiseev, Simon, Lvovsky, "Far-field linear optical superresolution via heterodyne detection in a higher-order local oscillator mode," Optica 3, 1148 (2016).
- Tham, Ferretti, Steinberg, "Beating Rayleigh's Curse by Imaging Using Phase Information," Phys. Rev. Lett. **118**, 070801 (2017).
- Paúr, Stoklasa, Hradil, Sánchez-Soto, Rehacek, "Achieving the ultimate optical resolution," Optica 3, 1144 (2016).



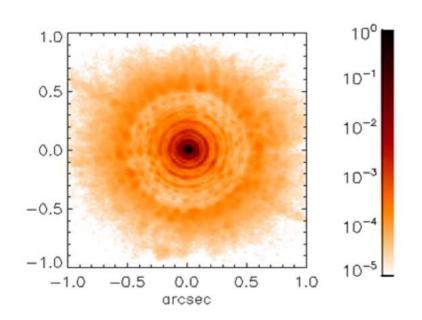
Imaging of One Point Source

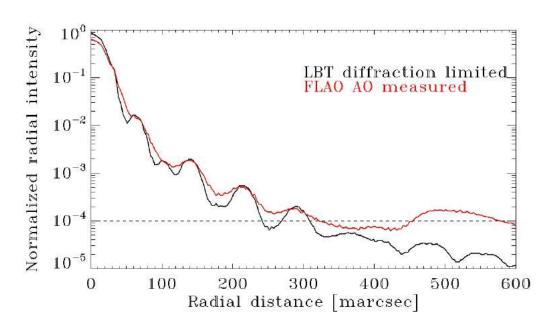


Diffraction limited

- 1. Fluorescence microscopy
- 2. Space telescopes (Webb, \$8.8 billion)
- 3. Ground-based telescopes:
 - (a) Large Binocular Telescope (LBT) (\$120 million)
 - (b) Giant Magellan Telescope (GMT) (\sim \$1 billion)
 - (c) Thirty Meter Telescope (TMT) (\sim \$1.2 billion)
 - (d) European Extremely Large Telescope (E-ELT) (\sim \$1.2 billion)

LBT Point-Spread Function





- Esposito et al., "Large Binocular Telescope Adaptive Optics System: New achievements and perspectives in adaptive optics," Proc. SPIE **8149**, 814902 (2011)
- Strehl ratio > 80% at infrared

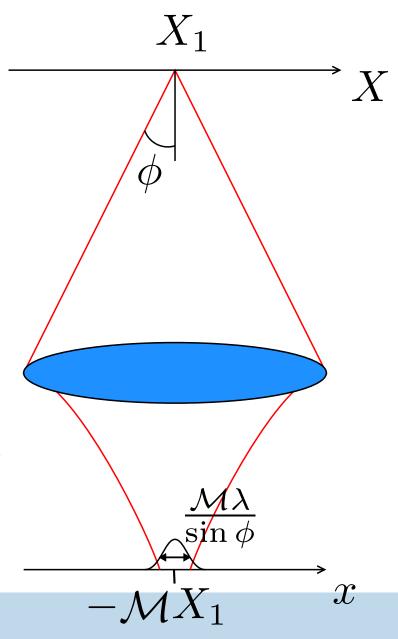
Inferring Position of One Point Source

- Classical source
- \blacksquare Given N detected photons, mean-square error:

$$\Sigma = \frac{\sigma^2}{N},\tag{1}$$

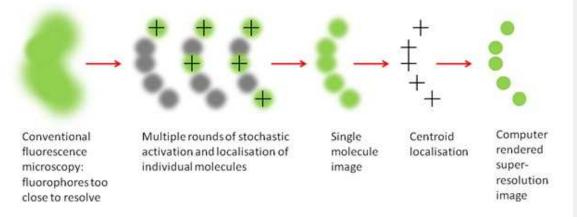
$$\sigma \sim \frac{\lambda}{\sin \phi}$$
. (2)

- Optics: Farrell (1966)
- Quantum: Helstrom (1970)
- Astronomy: Lindegren (1978)
- Microscopy: Bobroff (1986)
- Full EM, full quantum: Tsang, Optica 2, 646 (2015).



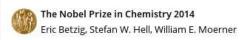
Superresolution Microscopy

■ PALM, STORM, etc.: **isolate** emitters. Locate **centroids**.



https://cam. facilities.northwestern. edu/588-2/single-molecule-localization-microscopy/

- Special fluorophores
- slow
- doesn't work for stars



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The Nobel Prize in Chemistry 2014







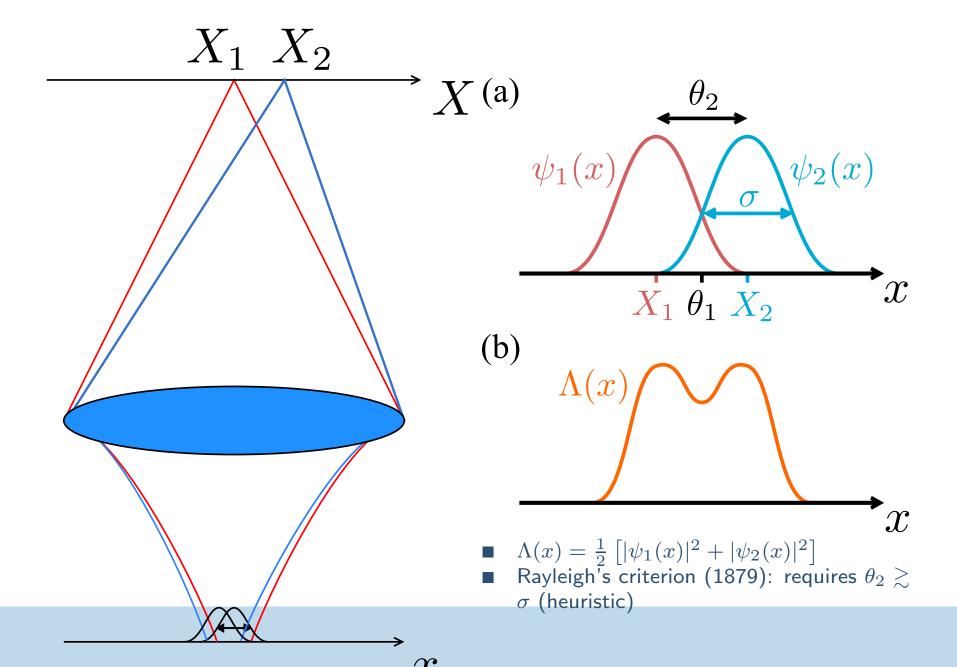
Photo: A. Mahmoud Stefan W. Hell Prize share: 1/3



Photo: A. Mahmoud William E. Moerner Prize share: 1/3

The Nobel Prize in Chemistry 2014 was awarded jointly to Eric Betzig, Stefan W. Hell and William E. Moerner "for the development of super-resolved fluorescence microscopy".

Two Point Sources



Centroid and Separation Estimation

- Bettens et al., Ultramicroscopy 77, 37 (1999); Van Aert et al., J. Struct. Biol. 138, 21 (2002); Ram, Ward, Ober, PNAS 103, 4457 (2006).
- Incoherent sources, Poisson statistics
- Cramér-Rao bound (standard in singlemolecule imaging)
- $X_1 = \theta_1 \theta_2/2, X_2 = \theta_1 + \theta_2/2.$
- centroid:

$$\Sigma_{11} \ge \frac{\sigma^2}{N}.\tag{3}$$

- CRB for **separation estimation**: two regimes
 - \bullet $\theta_2 \gg \sigma$:

$$\Sigma_{22} \ge \frac{4\sigma^2}{N},$$

 $\theta_2 \ll \sigma$:

$$\Sigma_{22} o \frac{4\sigma^2}{N} imes \infty$$
 (5)

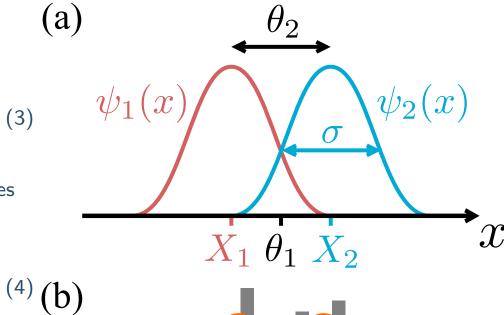
Beyond Rayleigh's criterion: A resolution measure with application to single-molecule microscopy

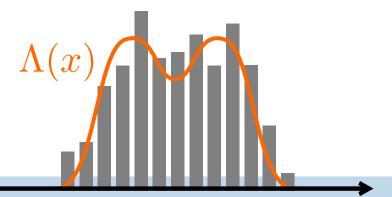
(a)

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Edited by Robert H. Austin, Princeton University, Princeton, NJ, and approved January 4, 2006 (received for review September 14, 2005)

Rayleigh's criterion is extensively used in optical microscopy for various experimental factors that affect the acquired data such





Rayleigh's Curse

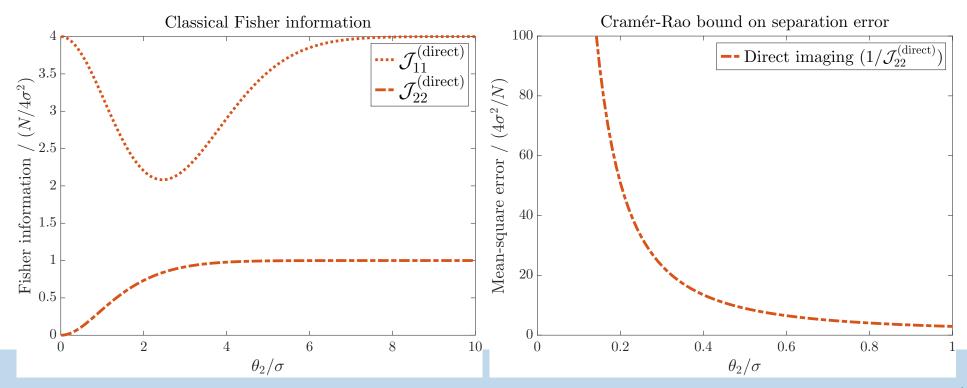
Cramér-Rao bound:

$$\Sigma_{11} \ge \frac{1}{\mathcal{J}_{11}^{\text{(direct)}}}$$

$$\Sigma_{22} \ge \frac{1}{\mathcal{J}_{22}^{\text{(direct)}}}$$
(6)

 $\mathcal{J}^{(\mathrm{direct})}$ is Fisher information for direct imaging

- Gaussian PSF, similar behavior for other PSF
- Rayleigh's curse
- PALM/STED/STORM: avoid violating Rayleigh



Quantum Information

- Measuring the spatial intensity (direct imaging, CCD) is just one measurement method. Quantum mechanics allows infinite possibilities.
- Helstrom (1967): For any measurement

$$\Sigma \ge \mathcal{J}^{-1} \ge \mathcal{K}^{-1}, \tag{7}$$

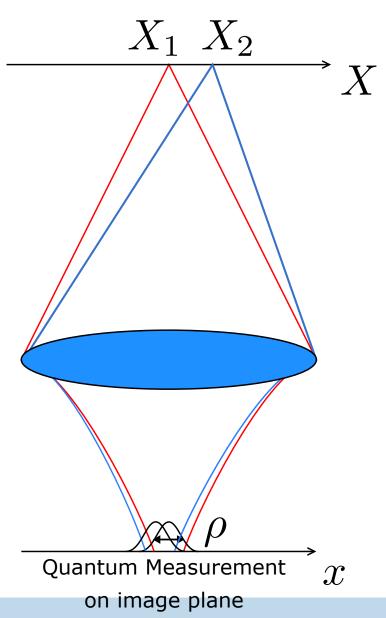
$$\mathcal{K}_{\mu\nu} = M \operatorname{Re} \left(\operatorname{tr} \mathcal{L}_{\mu} \mathcal{L}_{\nu} \rho \right), \tag{8}$$

$$\frac{\partial \rho}{\partial \theta_{\mu}} = \frac{1}{2} \left(\mathcal{L}_{\mu} \rho + \rho \mathcal{L}_{\mu} \right). \tag{9}$$

- $\mathcal{K}(\rho)$ is the quantum Fisher information, the ultimate amount of information in the photons.
- Coherent sources: Tsang, Optica 2, 646 (2015).
- Mixed states:

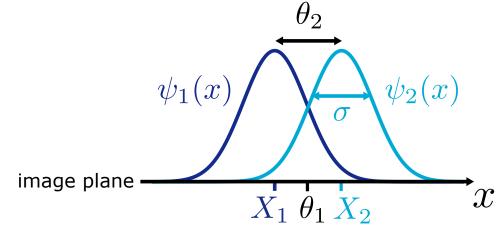
$$\rho = \sum_{n} D_n |e_n\rangle \langle e_n|, \qquad (10)$$

$$\mathcal{L}_{\mu} = 2 \sum_{n,m;D_n + D_m \neq 0} \frac{\langle e_n | \frac{\partial \rho}{\partial \theta_{\mu}} | e_m \rangle}{D_n + D_m} |e_n\rangle \langle e_m|. \quad (11)$$



Quantum Optics

- Mandel and Wolf, Optical Coherence and Quantum Optics; Goodman, Statistical Optics
- Thermal sources, e.g., stars, fluorescent particles.
- Average photon number per mode $\epsilon \ll 1$ at optical frequencies (visible, UV, X-ray, etc.).
- $\epsilon \sim 0.01$ for the sun at visible, $\epsilon \sim 10^{-6}$ for fluorophores.



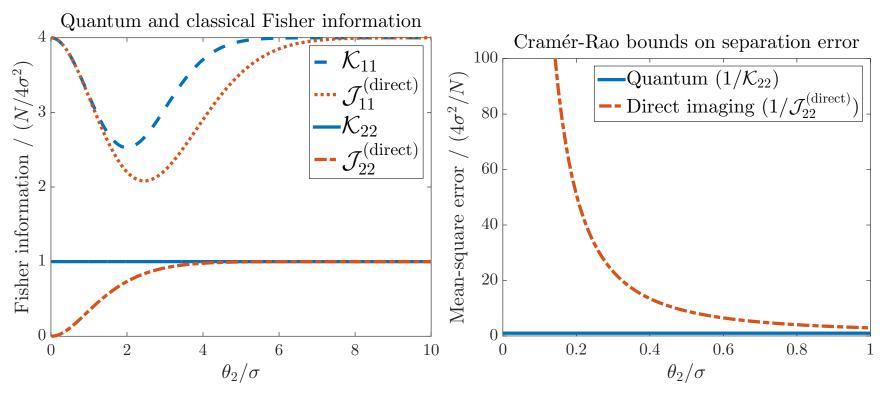
lacktriangle Quantum state in M temporal modes on **image plane** is $ho^{\otimes M}$, where

$$\rho = (1 - \epsilon) |\mathsf{vac}\rangle \langle \mathsf{vac}| + \frac{\epsilon}{2} (|\psi_1\rangle \langle \psi_1| + |\psi_2\rangle \langle \psi_2|) + O(\epsilon^2) \qquad [\langle \psi_1 | \psi_2 \rangle \neq 0,]$$
 (12)

$$|\psi_1\rangle = \int_{-\infty}^{\infty} dx \psi(x - X_1) |x\rangle, \quad |\psi_2\rangle = \int_{-\infty}^{\infty} dx \psi(x - X_2) |x\rangle.$$
 (13)

- derive from zero-mean Gaussian P function, mutual coherence
- Multiphoton coincidence: rare, little info as $\epsilon \ll 1$ (homeopathy)
- Similar model for stellar interferometry in Gottesman, Jennewein, Croke, PRL **109**, 070503 (2012); Tsang, PRL **107**, 270402 (2011).

Plenty of Room at the Bottom



■ Tsang, Nair, and Lu, Physical Review X 6, 031033 (2016)

$$\Sigma_{22} \ge \frac{1}{\mathcal{K}_{22}} = \frac{1}{N\Delta k^2}.\tag{14}$$

- **thermal sources with arbitrary** ϵ : Nair and Tsang, PRL (Editors' Suggestion) **117**, 190801 (2016); Lupo and Pirandola, *ibid.* **117**, 190802 (2016).
- Hayashi ed., Asymptotic Theory of Quantum Statistical Inference; Fujiwara JPA 39, 12489 (2006): there exists a POVM such that $\Sigma_{\mu\mu} \to 1/\mathcal{K}_{\mu\mu}$, $N \to \infty$.

Hermite-Gaussian Basis

project in Hermite-Gaussian basis:

$$E_1(q) = |\phi_q\rangle \langle \phi_q|, \qquad (15)$$

$$|\phi_q\rangle = \int_{-\infty}^{\infty} dx \phi_q(x) |x\rangle,$$
 (16)

$$\phi_q(x) = \left(\frac{1}{2\pi\sigma^2}\right)^{1/4} H_q\left(\frac{x}{\sqrt{2}\sigma}\right) \exp\left(-\frac{x^2}{4\sigma^2}\right). \tag{17}$$

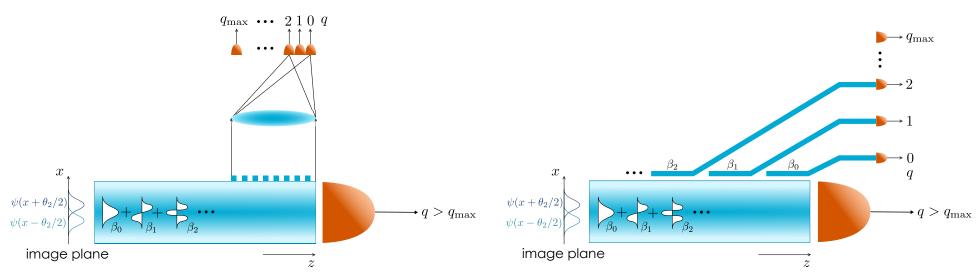
Assume PSF $\psi(x)$ is Gaussian (common).

$$\frac{1}{\mathcal{J}_{22}^{(\text{HG})}} = \frac{1}{\mathcal{K}_{22}} = \frac{4\sigma^2}{N}.$$
 (18)

- \blacksquare Maximum-likelihood estimator can saturate the classical bound asymptotically for large N.
- arXiv:1605.03799v2: SPADE with Max-Like:

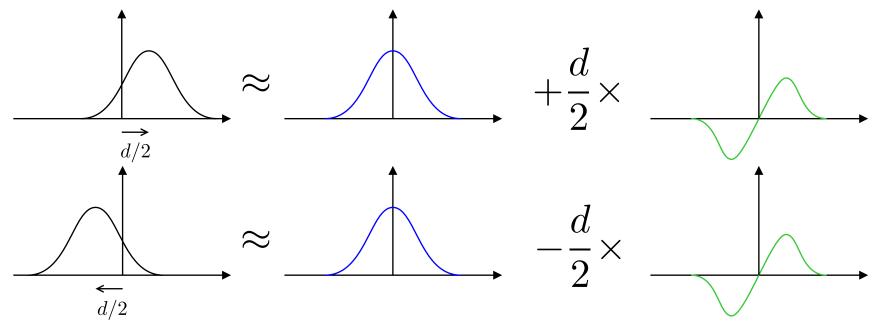
$$\Sigma_{22} \le \frac{16\sigma^2}{N}$$
 for any detected N . (19)

Spatial-Mode Demultiplexing (SPADE)



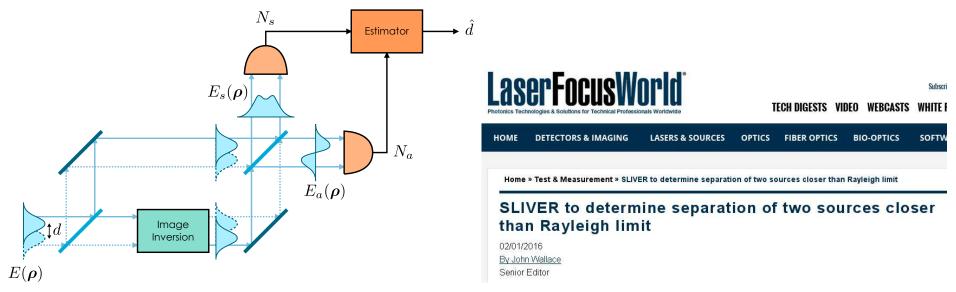
- Many other ways (optical comm.), e.g.,
 - ◆ DAB Miller, "Self-configuring universal linear optical components," Photonics Research 1, 1 (2013).
 - ◆ Guifang Li et al., "Space-division multiplexing: the next frontier in optical communication," Adv. Opt. Photon. 6, 413 (2014).
 - ◆ V. A. Soifer, Computer Design of Diffractive Optics (CISP/Woodhead, Cambridge, 2013)

Elementary Explanation



- Incoherent sources: energy in first-order mode is $\propto (d/2)^2 + (-d/2)^2 = d^2/2$
- Zeroth-order mode is just background noise, removing it improves SNR.
- Why quantum formalism?
 - ◆ Fundamental quantum limit
 - ♦ Ensures measurement is physical
 - Discover new possibilities

SLIVER



- SuperLocalization via Image-inVERsion interferometry
- Nair and Tsang, Opt. Express **24**, 3684 (2016).
- Laser Focus World, Feb 2016 issue.
- Classical theory/experiment of image-inversion interferometer:
 - ◆ Wicker, Heintzmann, Opt. Express 15, 12206 (2007).
 - ♦ Wicker, Sindbert, Heintzmann, Opt. Express 17, 15491 (2009).
 - ◆ Weigel, Babovsky, Kiessling, Kowarschik, Opt. Commun. 284, 2273 (2011).
 - ◆ No statistical analysis, predicted <2X resolution enhancement

Theoretical Follow-up

		Dimensions	Sources	Theory	Experimental proposals
1.	Tsang, Nair, and Lu, Phys. Rev. X 6 , 031033 (2016)	1D	Weak thermal (optical frequencies and above)	Quantum	SPADE
2.	Nair and Tsang, Optics Express 24 , 3684 (2016)	2D	Thermal (any fre- quency)	Semiclassical	SLIVER
3.	Tsang, Nair, and Lu, Proc. SPIE 10029 , 1002903 (2016)	N/A	Weak thermal, lasers	Semiclassical	N/A
4.	Nair and Tsang, PRL (Editors' Suggestion) 117, 190801 (2016)	1D	Thermal	Quantum	SLIVER
5.	Tsang, arXiv:1605.03799	1D	Weak thermal	Quantum, Bayesian, Minimax	SPADE
6.	Ang, Nair, and Tsang, arXiv:1606.00603	2D	Weak thermal	Quantum	SPADE, SLIVER
7.	Tsang, arXiv:1608.03211	2D	Weak thermal, multiple sources	Quantum	SPADE
8.	Lu, Nair, Tsang, arXiv:1609.03025	2D	Weak thermal, one - versus-two	Quantum, binary hy- pothesis testing	SPADE, SLIVER

Other groups:

- Lupo and Pirandola, PRL (Editors' Suggestion) 117, 190801 (2016).
- Rehacek *et al.*, Optics Letters **42**, 231 (2017).
- Krovi, Guha, Shapiro, arXiv:1609.00684.
- Kerviche, Guha, Ashok, arXiv:1701,04913.

Experiments

- Tang, Durak, and Ling, "Fault-tolerant and finite-error localization for point emitters within the diffraction limit," Optics Express **24**, 22004 (2016).
 - SLIVER
 - ♦ Laser, classical noise
- Yang, Taschilina, Moiseev, Simon, Lvovsky, "Far-field linear optical superresolution via heterodyne detection in a higher-order local oscillator mode," Optica 3, 1148 (2016).
 - Mode heterodyne
 - ◆ Laser
- Tham, Ferretti, Steinberg, "Beating Rayleigh's Curse by Imaging Using Phase Information," Phys. Rev. Lett. 118, 070801 (2017).
 - variation of SPADE
 - two independent SPDC sources, $\sim 2 \times$ quantum limit
- Paúr, Stoklasa, Hradil, Sánchez-Soto, Rehacek, "Achieving the ultimate optical resolution," Optica 3, 1144 (2016).
 - variation of SPADE
 - laser, close to quantum limit (caveat)

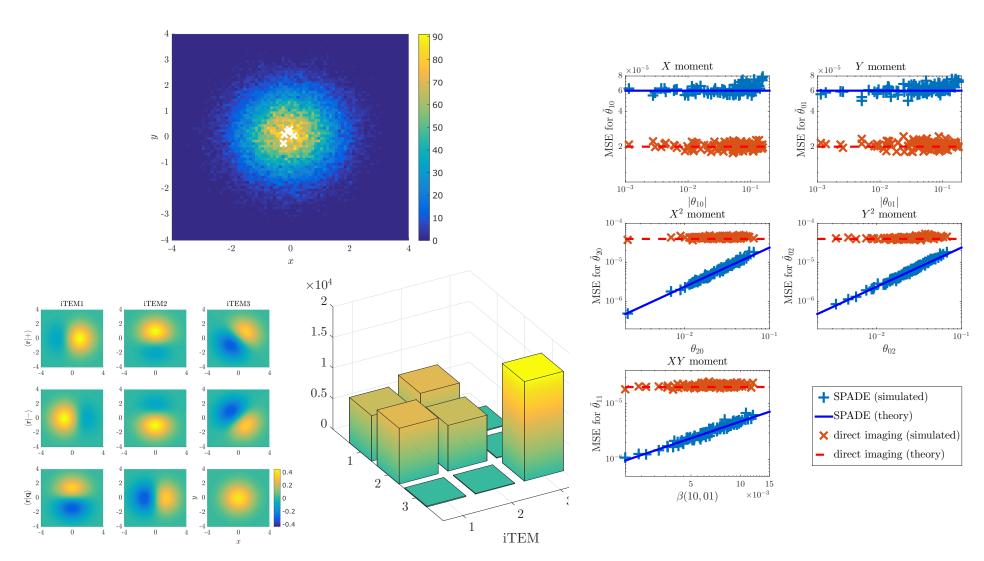
SPLICE

■ Tham, Ferretti, Steinberg, Phys. Rev. Lett. 118, 070801 (2017).

image plane

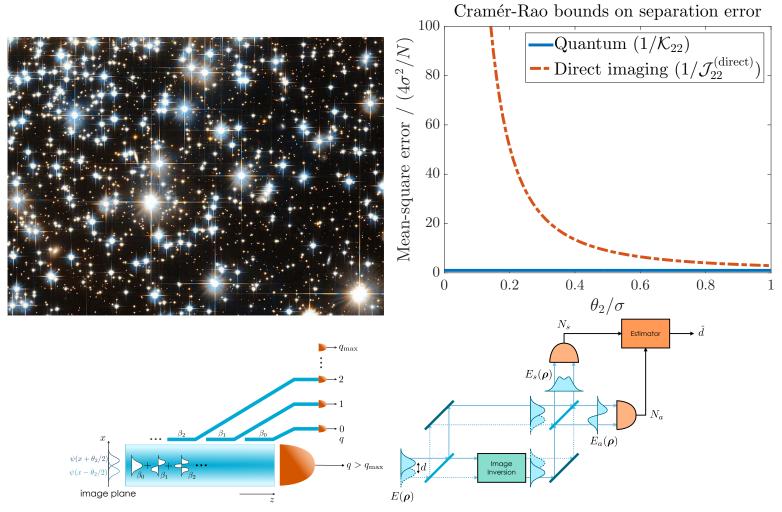
| Single-mode fiber | SPLICE, fire (N=1213) | SPLICE, coarse (N=1451) | O (PC, fire, 2nd ord, in reg. (N=2792) | O (PC, fire, 2nd ord, in reg. (N=2890) | O (PC, coarse, built-in (N=2890) | O (PC, coarse, built-in

Arbitrary Source Distributions



- Yang *et al.*, Optica **3**, 1148 (2016): even moments
- Tsang, arXiv:1608.03211: Generalized SPADE: Enhanced estimation of 2nd or higher moments

Quantum Metrology Kills Rayleigh's Criterion



- FAQ: https://sites.google.com/site/mankeitsang/news/rayleigh/faq
- email: mankei@nus.edu.sg

$\epsilon \ll 1$ Approximation

■ Chap. 9, Goodman, *Statistical Optics*:

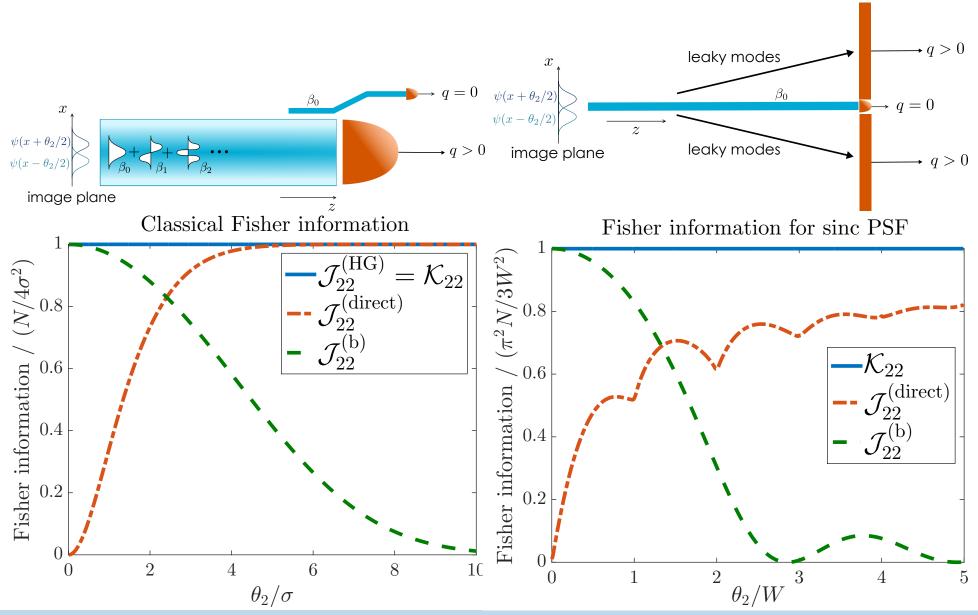
"If the count degeneracy parameter is much less than 1, it is highly probable that there will be either zero or one counts in each separate coherence interval of the incident classical wave. In such a case the classical intensity fluctuations have a negligible "bunching" effect on the photo-events, for (with high probability) the light is simply too weak to generate multiple events in a single coherence cell.

Zmuidzinas (https://pma.caltech.edu/content/jonas-zmuidzinas), JOSA A 20, 218 (2003):

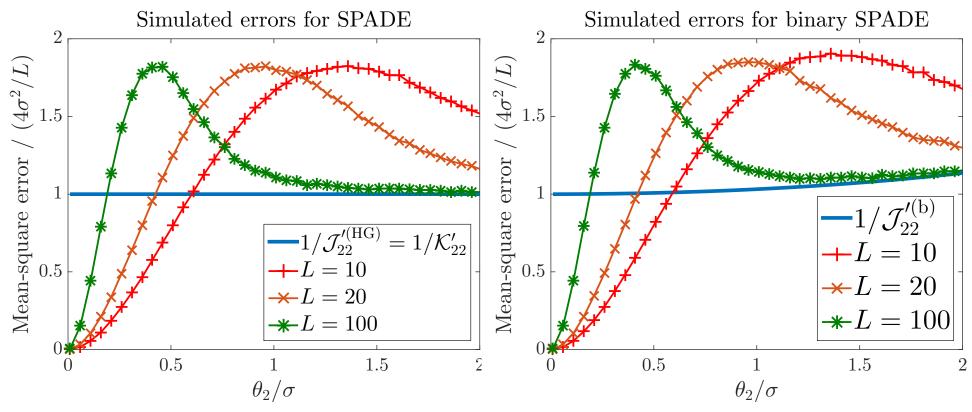
"It is well established that the photon counts registered by the detectors in an optical instrument follow statistically independent Poisson distributions, so that the fluctuations of the counts in different detectors are uncorrelated. To be more precise, this situation holds for the case of thermal emission (from the source, the atmosphere, the telescope, etc.) in which the mean photon occupation numbers of the modes incident on the detectors are low, $n \ll 1$. In the high occupancy limit, $n \gg 1$, photon bunching becomes important in that it changes the counting statistics and can introduce correlations among the detectors. We will discuss only the first case, $n \ll 1$, which applies to most astronomical observations at optical and infrared wavelengths."

- Hanbury Brown-Twiss (post-selects on two-photon coincidence, homeopathy): poor SNR, obsolete for decades in astronomy.
- See also Labeyrie et al., An Introduction to Optical Stellar Interferometry, etc.
- Fluorescent particles: Pawley ed., Handbook of Biological Confocal Microscopy, Ram, Ober, Ward (2006), etc., may have antibunching, but Poisson model is fine and standard because of $\epsilon \ll 1$.

Binary SPADE



Numerical Performance of Maximum-Likelihood Estimators

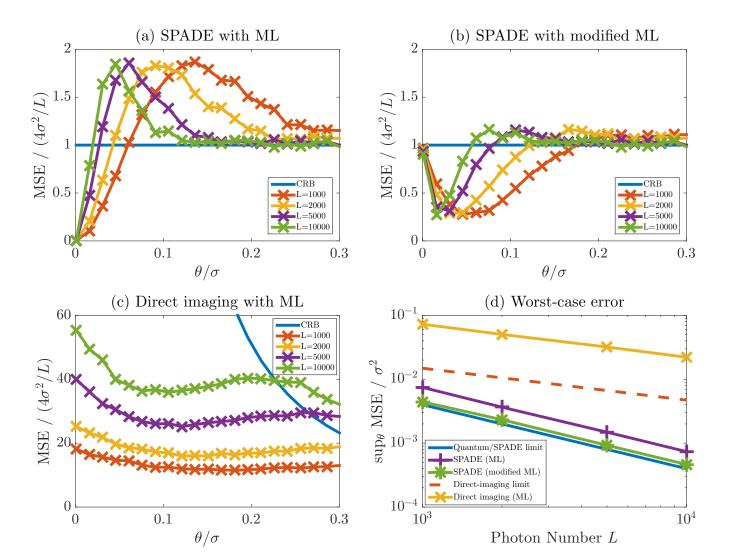


- lacksquare L= number of detected photons
- **biased** (violate CRB), $< 2 \times$ CRB.

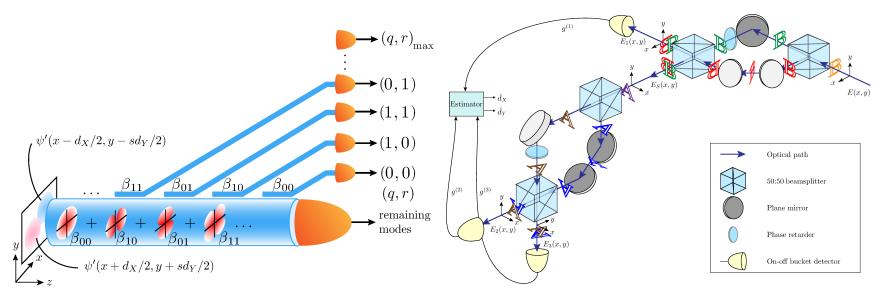
Minimax/Bayesian

Bayesian CRB for any biased/unbiased estimator (e-print arXiv:1605.03799)

Quantum/SPADE:
$$\sup_{\theta} \Sigma_{22}(\theta) \ge \frac{4\sigma^2}{N}$$
, Direct imaging: $\sup_{\theta} \Sigma_{22}^{(\text{direct})}(\theta) \ge \frac{\sqrt{2}\sigma^2}{3\sqrt{N}}$. (20)

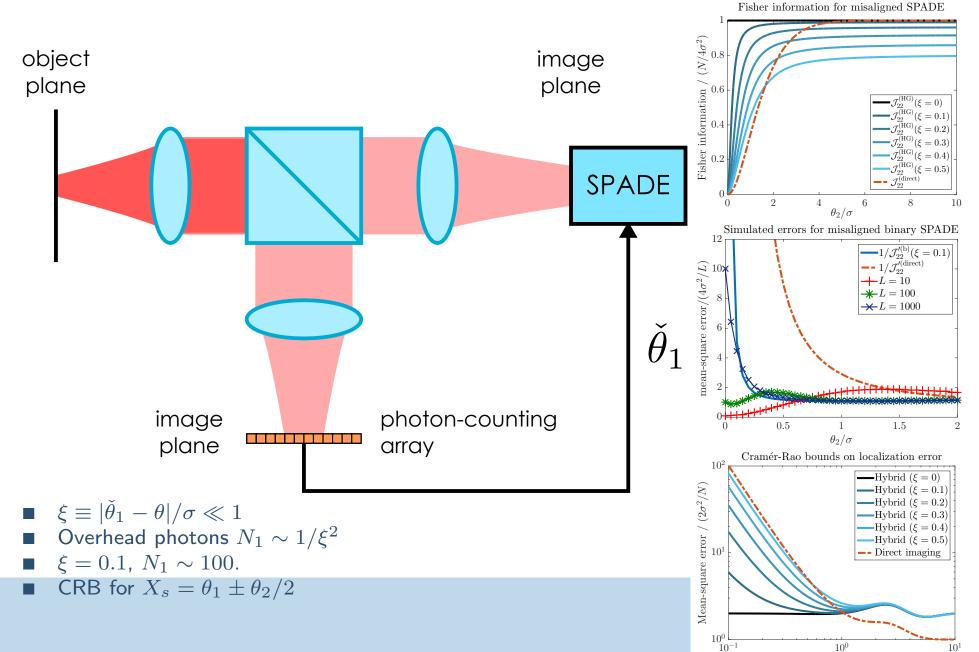


2D SPADE and SLIVER

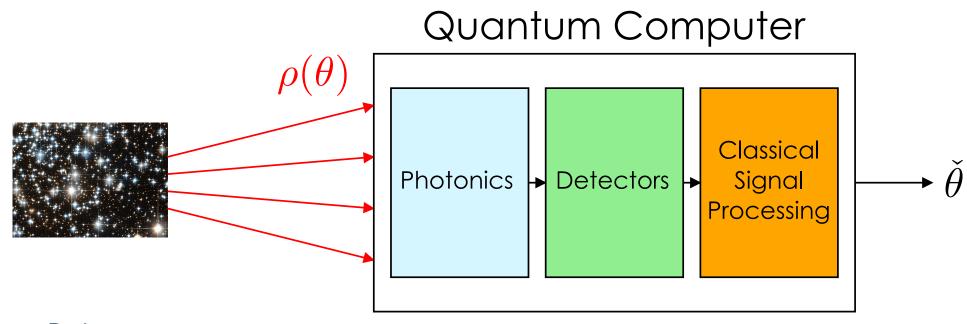


Ang, Nair, Tsang, e-print arXiv:1606.00603

Misalignment



Quantum Computational Imaging



- Design quantum computer to
 - Maximize information extraction
 - Reduce classical computational complexity

Quantum Technology 1.5

