Supporting Information

Nanometer-Scale Heterogeneous Interfacial Sapphire Wafer-Bonding for Enabling Plasmonic-Enhanced Nanofluidic Mid-Infrared Spectroscopy

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Note S1: Simulated electrical field distribution



Figure S1. Simulated electrical field distribution along Z direction with Au reflector at different nanogaps: (a) nanogap = 150 nm; (b) nanogap = 500 nm. The quadrupole resonance mode is easily affected by the gap distance. Not only the quadrupole resonance will disappear, but also the resonance intensity will also be weakened when the nanogap is too large.



Note S2: Extraction of external and intrinsic losses

Figure S2. Fitting results of reflection spectra at different nanogaps according to the temporal coupled-mode theory (TCMT).

Gap	22	30	33	44	55	66	77	88	90
γr	0.9013	1.1441	1.1788	1.4180	1.6487	1.8597	2.0815	2.3560	2.3947
γa	4.2836	3.9208	3.8018	3.4522	3.1581	2.8315	2.6460	2.4500	2.3024
γ_r/γ_a	0.2104	0.2918	0.3100	0.4108	0.5220	0.6568	0.7867	0.9616	1.0401
\mathbf{R}_1	42.45%	29.01%	27.76%	17.27%	9.89%	4.87%	1.50%	1.10%	1.51%
\mathbf{R}_2	82.11%	77.25%	75.60%	69.08%	63.11%	56.93%	50.63%	44.73%	43.01%
$\gamma_r/(\gamma_a+\gamma_m)$	0.0494	0.0598	0.0696	0.0926	0.1144	0.1389	0.1688	0.1890	0.1961
Gap	99	110	121	132	140	143	154	165	176
γr	2.5979	3.0686	3.3873	3.7650	3.9971	4.0620	4.3813	4.6999	4.9621
γ_{a}	2.1979	1.9405	1.7951	1.7020	1.6580	1.6273	1.6101	1.5769	1.5431
γ_r/γ_a	1.1820	1.5813	1.8870	2.2121	2.4108	2.4962	2.7211	2.9805	3.2158
\mathbf{R}_1	1.90%	5.08%	10.79%	15.04%	17.65%	18.48%	21.82%	24.93%	28.04%
\mathbf{R}_2	37.81%	32.88%	28.70%	24.66%	21.97%	21.50%	18.63%	16.30%	13.98%
$\gamma_r/(\gamma_a+\gamma_m)$	0.2180	0.2640	0.2980	0.3376	0.3483	0.3689	0.3984	0.4254	0.4576

Table S1. Extraction of external and intrinsic losses at different nanogaps based on TCMT.



Figure S3. Relationship between γ_r/γ_a and $\gamma_r/(\gamma_a + \gamma_m)$ according to the extracted data.

Note S3: Theoretical calculation of the signal enhancement

In section S3, we demonstrate the theoretical calculation of the signal enhancement based on the temporal coupled-mode theory (TCMT). We can derive the signal enhancement by the difference between the spectra with and without the analyte absorption. Detailed theoretical calculations are as follows:

(1) Without analyte absorption:

General equations of the TCMT can be described as follows:

$$\frac{d\mathbf{P}}{dt} = jw_0P - (\gamma_a + \gamma_r)P + ks_+$$
(S1)

$$s_{-} = -s_{+} + k\mathbf{M} \tag{S2}$$

where *P* is the mode amplitude of the plasmonic structure. ω_0 is the resonance frequency of the plasmonic structure. We can obtain the following formulas by calculating (S1) and (S2):

$$jwA = jw_0A - (\gamma_a + \gamma_r)A + ks_+$$
(S3)

$$s_{+} = \frac{j(w - w_0) + (\gamma_a + \gamma_r)}{k} A$$
 (S4)

$$s_{-} = \frac{k^2 - j(w - w_0) - (\gamma_a + \gamma_r)}{k} A$$
(S5)

Therefore, the calculated results of the reflection spectra from (S1) and (S2) are as follows:

$$R = \left|\frac{s_{-}}{s_{+}}\right|^{2} = \left|\frac{k^{2} - j(w - w_{0}) - (\gamma_{a} + \gamma_{r})}{j(w - w_{0}) + (\gamma_{a} + \gamma_{r})}\right|^{2}$$
(S6)

Due to the plasmonic nanofluidic is single port resonance system, so:

$$k = \sqrt{2\gamma_{\rm r}}$$

Changes in reflection spectra with frequency (w) can be expressed as follows:

$$R = \left| \frac{s}{s}_{+} \right|^{2} = \left| \frac{j(w - w_{0}) - (\gamma_{r} - \gamma_{a})}{j(w - w_{0}) + (\gamma_{a} + \gamma_{r})} \right|^{2}$$
(S7)

The reflection dip can be obtained when $w = w_0$:

$$R = \left| \frac{\gamma r - \gamma_a}{\gamma_a + \gamma_r} \right|^2$$
(S8)

(2) With analyte absorption:

By adding the analyte absorption, the derivations are as follows:

$$\frac{dP}{dt} = jw_0P - (\gamma_a + \gamma_r)A + j\mu M + ks_+$$
(S9)

$$\frac{dM}{dt} = jw_m M - \gamma_m M + j\mu P \tag{S10}$$

$$s_{-} = -s_{+} + kA \tag{S11}$$

$$R = \left| \frac{s_{-}}{s_{+}} \right|^2 \tag{S12}$$

$$s_{+} = \frac{j(w - w_0) + (\gamma_a + \gamma_r) + \frac{\mu^2}{j(w - w_m) + \gamma_m}}{k} P$$
(S13)

$$s_{-} = \frac{k^2 - j(w - w_0) - (\gamma_a + \gamma_r) - \frac{\mu^2}{j(w - w_m) + \gamma_m}}{k}P$$
(S14)

$$R_{A} = \left|\frac{s}{s}\right|^{2} = \left|\frac{k^{2} - j(w - w_{0}) - (\gamma_{a} + \gamma_{r}) - \frac{\mu^{2}}{j(w - w_{m}) + \gamma_{m}}}{j(w - w_{0}) + (\gamma_{a} + \gamma_{r}) + \frac{\mu^{2}}{j(w - w_{m}) + \gamma_{m}}}\right|^{2}$$
(S15)

where μ is the coupling strength. *M* is the molecular vibration. ω_m is the resonance frequency of the molecular vibration. When $w = w_0$,

$$R_{A} = \left| \frac{(\gamma_{r} - \gamma_{a}) - \frac{\mu^{2}}{j(w - w_{m}) + \gamma_{m}}}{(\gamma_{a} + \gamma_{r}) + \frac{\mu^{2}}{j(w - w_{m}) + \gamma_{m}}} \right|^{2} = \left| \frac{\gamma_{r} - \left(\gamma_{a} + \frac{\mu^{2}}{j(w - w_{m}) + \gamma_{m}}\right)}{\gamma_{r} + \left(\gamma_{a} + \frac{\mu^{2}}{j(w - w_{m}) + \gamma_{m}}\right)} \right|^{2}$$
(S16)

We assume that:

$$\gamma_{\mu} = \frac{\mu^2}{j(w - w_m) + \gamma_m} \tag{S17}$$

$$\gamma'_a \rightarrow \gamma_a + \gamma_m \tag{S18}$$

Therefore, we can get the reflection by submitting (S17) and (S18) into (S16):

$$R_A = \left| \frac{\gamma_r - \gamma'_a}{\gamma_r + \gamma'_a} \right|^2 \tag{S19}$$

(3) Enhanced signal:

The enhanced signal is the difference between the coupling with and without analyte absorptions at the resonance frequency:

$$\Delta R_A = R_{H_20} - R = \left| \frac{\gamma_r - \gamma'_a}{\gamma_r + \gamma'_a} \right|^2 - \left| \frac{\gamma_r - \gamma_a}{\gamma_r + \gamma_a} \right|^2$$
(S20)

We assume that:

$$f = \frac{\gamma_r}{\gamma_a} \tag{S21}$$

$$f' = \frac{\gamma'_r}{\gamma_a} \tag{S22}$$

Then we can get the enhanced signal:

$$\Delta R_A = \left| \frac{f' - 1}{f' + 1} \right|^2 - \left| \frac{f - 1}{f + 1} \right|^2 = 4 \frac{(f'f - 1)(f' - f)}{(f' + 1)^2(f + 1)^2}$$
(S23)

According to the result fitting result that f' is approximately positively correlated to f:

$$f' \cong tf$$
 (S24)

Therefore, the enhanced signal can be obtained:

$$\Delta R_A \propto \frac{(tf^2 - 1)(t - 1)f}{(tf + 1)^2(f + 1)}$$
(S25)

Note S4: Bonding strength measurement



Figure S4. Measurement for the Al_2O_3/SiO_2 bonding energy. (a) The schematic diagram for the bonding strength calculation. (b) Optical image of the bonded Al_2O_3/SiO_2 wafer with the blade insertion.



Note S5: Surface characterization of SiO₂ and Al₂O₃ surfaces by XPS

Figure S5. Surface characterization by XPS: (a, b) Si 2p and O 1s core spectra for SiO₂ substrates before and after the plasma activation. (c, d) Al 2p and O 1s core spectra for Al₂O₃ substrates before and after the plasma activation.





Figure S6. Spectra tuning and enhanced H_2O signal detection at different nanogaps: Measured reflection spectra of the plasmonics-nanofluidics filled with air and DI water at the gap of (a) 30 nm, (b) 90 nm, (c) 140 nm, and (d) 165 nm. To tune the resonance spectra, the nanoantenna length is varied from 572 nm to 668 nm with an interval of 12 nm.





Figure S7. Measurement for nanogaps. (a) AFM image for Au antennas fabricated on the Al_2O_3 substrate. The curve below the AFM image shows the antenna profiles along the height direction. The antenna heights are around 45 nm. (b)-(f) Etching profiles of SiO₂ substrates for 30 nm, 55 nm, 90 nm, 140 nm, and 165 nm-nanogap plasmonic nanofluidics.

Note S8: Device fabrication process



Figure S8. The fabrication process of plasmonics-nanofluidics device: (a-e) fabrication of the sensing chamber with Au reflector. To avoid the deposited Au film will cause the protrusion of the edges and induce the bonding failure, we repeat the laser writer and SiO₂ etching processes twice to avoid the contact between the deposited Au film and edges. (f-h) nanoantenna fabrication on Al_2O_3 substrate by electron beam lithography and electron beam evaporator. (i-j) plasma-activated direct bonding of Al_2O_3 and SiO₂ to integrate the plasmonics-nanofluidic device. The insert in the right corner is the modified plasma-activated direct bonding equipment.

Note S9: Calculation for the number of acetone molecules

According to the simulations of the electric field distribution shown in Figure 1e, the sensing area is mainly concentrated in the space sandwiched by the nanoantennas and Au reflector. The schematic diagram of the plasmonic nanofluidic is shown in Figure S8.



Figure S9. Schematic diagram of the plasmonic nanofluidic. (a) Au nanoantenna array fabricated on the sapphire substrate. The length (*L*) of the Au nanoantenna array is 200 μ m. (b) One unit of the Au nanoantenna array. The period (*P*) is 1.5 μ m. The length (*l*) and width (*w*) of the Au nanoantenna are 608 nm and 200 nm, respectively. (c) Cross-section of the plasmonic nanofluidic. The nanogap (*d*) between the Au nanoantenna and SiO₂ substrate is 55 nm.

Therefore, the total sensing volume can be expressed as follows:

$$V_T = L^2 \cdot \frac{l \cdot w}{P^2} \cdot d, \tag{S26}$$

where V_T is the total sensing volume, L is the length of the Au nanoantenna array, l is the nanoantenna length, w is the nanoantenna width, P is the period of the nanoantenna array, and d is nanogap between the nanoantenna and SiO₂ substrate.

So the number of acetone molecules can be obtained as follows:

$$N = \frac{V_T \cdot C_A}{V_A},\tag{S27}$$

where C_A and V_A are the concentration and molar volume of acetone, respectively. The V_A is

 $75.1 \times 10^{-9} \text{ m}^{3/\text{mol.}}$

By submitting (S26) into (S27), the final equation for the number of acetone molecules is obtained as follows:

$$N = L^2 \cdot \frac{l \cdot w}{P^2} \cdot d \cdot \frac{C_A}{75.1 \times 10^{-9}}.$$
 (S28)

Note S10: Dynamic monitoring for ethanol diffusion



Figure S10. Dynamic monitoring for ethanol diffusion. Measured reflection spectra when the plasmonic-enhanced nanofluidic filled with (a) acetone, (b) DI water, and (d) ethanol. Differential absorption spectra for DI water and ethanol are shown in (c) and (e), respectively. (f) Statistical results of absorbance, reflection dips, and reflection differences at different times.